

A New Technique for Solving Poisson's Equation with High Accuracy on Domains of Any Aspect Ratio

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ICFA Beam Dynamics Workshop on Space-Charge Simulation
Trinity College, Oxford
April 2-4, 2003

Performed in Collaboration with



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Critical Issue



- **Poisson solvers used in quasistatic electric and gravitational particle-in-cell simulations generally fail when the grid aspect ratio $\gg 1$**
- **Some important problems involve extreme aspect ratios:**
 - Long beams in rf accelerators: length $\sim 1\text{m}$; radius $\sim 1\text{mm}$
 - Beams in induction linacs: length $\sim 10\text{s of meters}$; radius $\sim \text{cm}$
 - Galaxies
- **Standard grid-based approaches involve using a very large # of grid points in the long dimension, leading to prohibitively long run times**
- **As a result, it is *extremely* difficult model high aspect ratio systems accurately using standard grid-based approaches**

Observations



- The Green function, G , and source density, ρ , may change over vastly different scales
- G is known apriori; ρ is not

We should use all the information available regarding G so that the numerical solution is only limited by our approximate knowledge of ρ

- Example: 2D Poisson equation in free space

$$\phi(x, y) = \iint G(x - x', y - y') \rho(x', y') dx' dy'$$

Standard Approach (Hockney and Eastwood)



$$\phi_{i,j} = \sum G_{i-i', j-j'} \phi_{i',j'}$$

- This approach is equivalent to using the trapezoidal rule (modulo treatment of boundary terms) to approximate the convolution integral
- This approach makes use of only partial knowledge of G
- The error depends on how rapidly the integrand, ϕG , varies over an elemental volume
 - If ϕ changes slowly we might try to use a large grid spacing; but this can introduce huge errors due to the change in G over a grid length

Cellular Analytic Convolution (CAC)



- Assume the charge density, ρ , varies in a prescribed way in each cell
- Use the analytic form of the Green function to perform the convolution integral exactly in each cell, then sum over cells
- Example: linear basis functions to approximate ρ in a cell:

$$\begin{aligned} \rho(x_i, y_j) = & \sum_{i', j'} \rho_{i', j'} \int_0^{h_x} dx' \int_0^{h_x} dy' (h_x - x')(h_y - y') G(x_i - x_i' - x', y_j - y_j' - y') + \\ & \sum_{i', j'} \rho_{i+1, j'} \int_0^{h_x} dx' \int_0^{h_x} dy' x' (h_y - y') G(x_i - x_i' - x', y_j - y_j' - y') + \\ & \sum_{i', j'} \rho_{i, j+1} \int_0^{h_x} dx' \int_0^{h_x} dy' (h_x - x') y' G(x_i - x_i' - x', y_j - y_j' - y') + \\ & \sum_{i', j'} \rho_{i+1, j+1} \int_0^{h_x} dx' \int_0^{h_x} dy' x' y' G(x_i - x_i' - x', y_j - y_j' - y') \end{aligned}$$

- Shifting the indices results in a single convolution involving an integrated effective Green function:

$$\rho_{i, j} = \sum_{i', j'} G_{i-i', j-j'}^{eff} \rho_{i', j'}$$

Cost and Accuracy of CAC; Improvement over Standard Approach



- **Cost:** Computing the elemental integrals can be done via analytical formulae in 2D (and 3D?) or by num. quadrature
 - Requires more FLOPS than simply using G_{ij} but...
 - In situations where the grid is fixed, this only needs to be done once at the start of a run. Amortized over many time steps, this does not significantly impact run time.
- **Accuracy:** Method works as long as the elemental integrals are computed accurately and as long as the grid and # of macroparticles are sufficient to resolve variation in \square
- CAC maintains accuracy even for extreme aspect ratios (>1000:1)

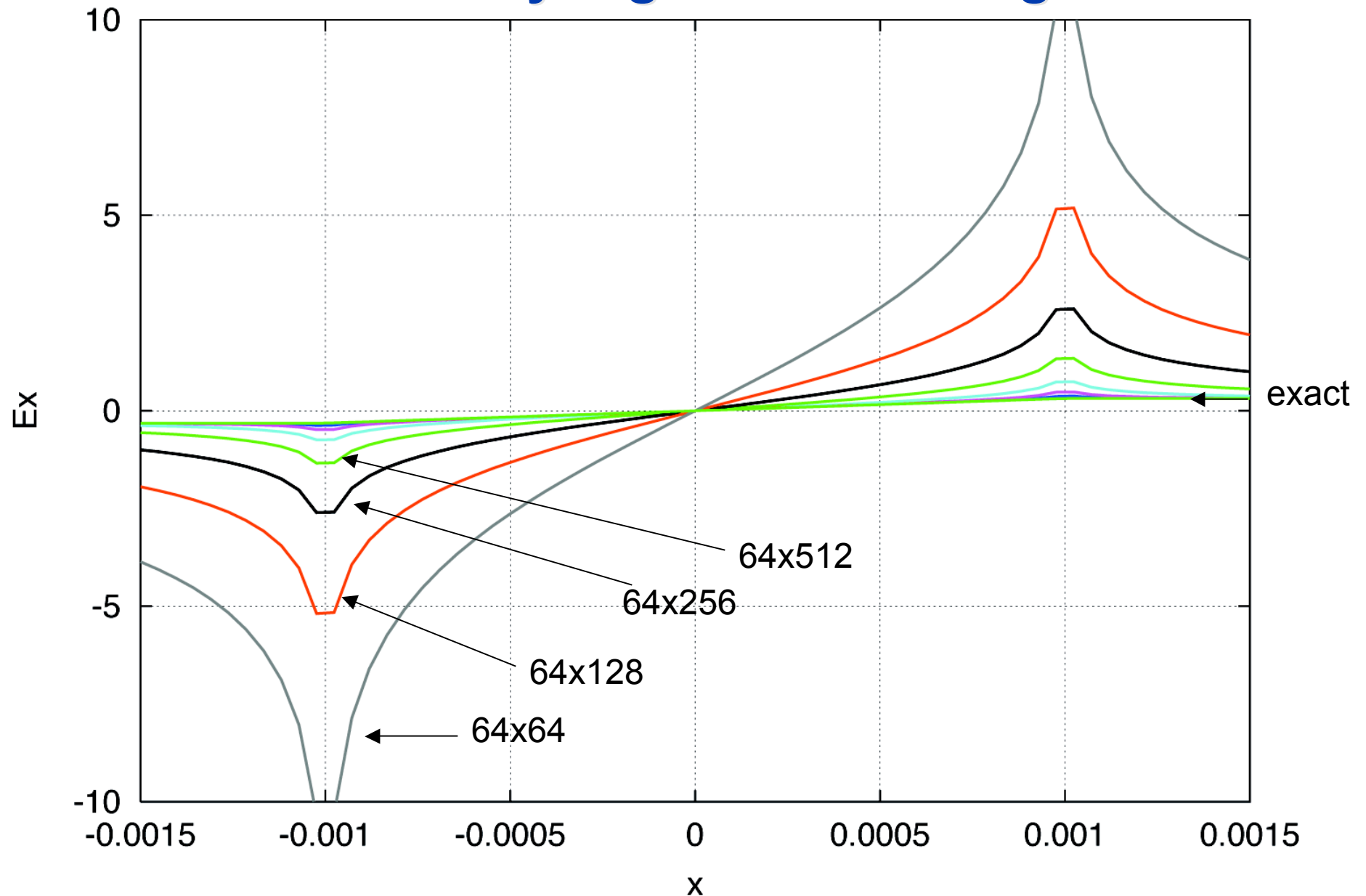
As a result, CAC performs orders of magnitude better than the standard convolution algorithm for realistic problems involving large aspect ratios

Example: Uniformly filled 2D ellipse

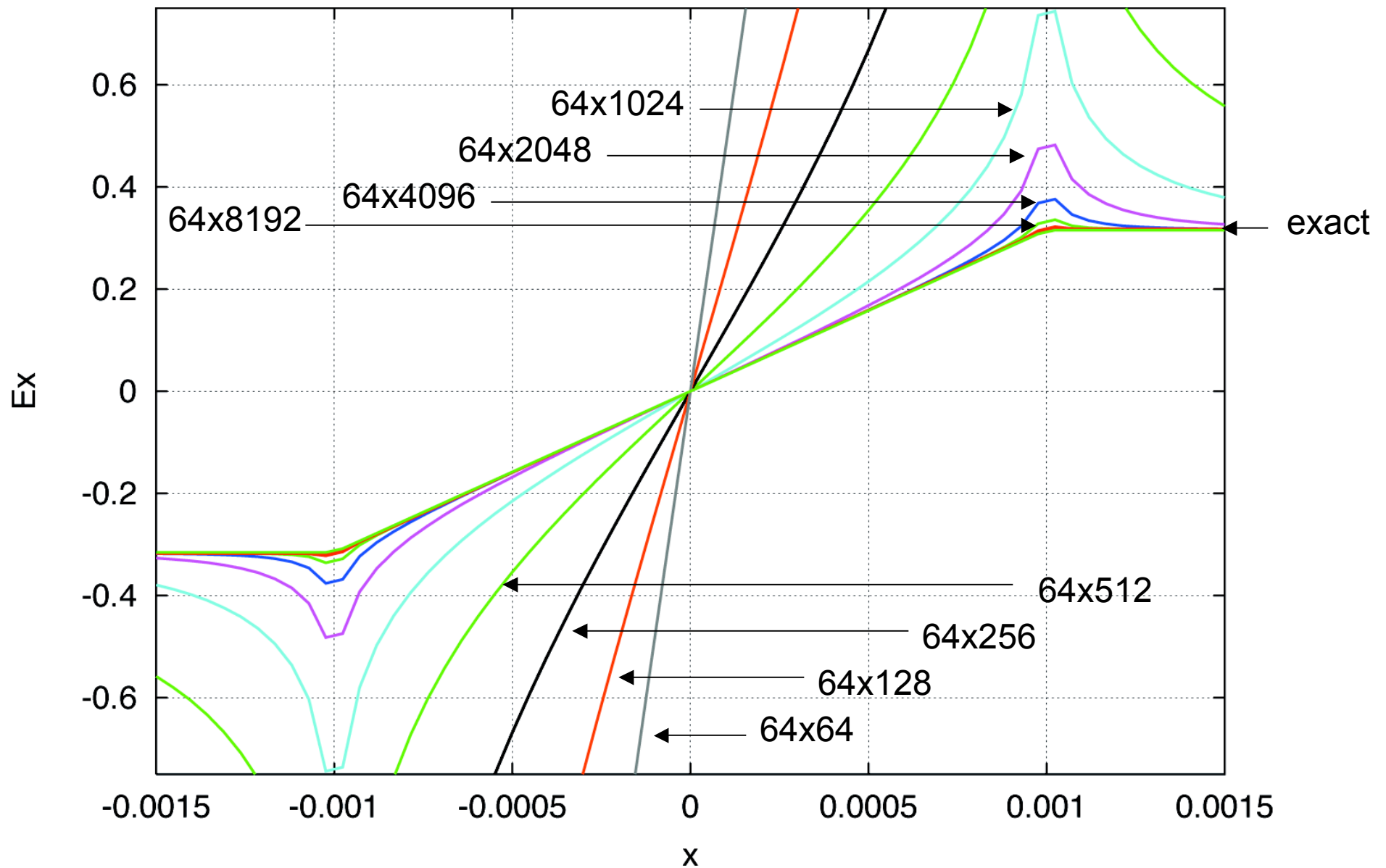


- **Aspect ratio is 1:1000 -- $x_{\max}=0.001$, $y_{\max}=1$**
- **Calculation of fields using (1) standard Hockney algorithm and (2) new CAC approach**
 - In both cases, performed convolutions for the fields directly (rather than calculating the potential and using finite differences to obtain fields)
- **Calculation performed on a grid of size $\pm 0.0015 \times \pm 1.5$ using a mesh of size**
 - Hockney: 64x64, 64x128, 64x256, ..., 64x16384
 - CAC: 64x64

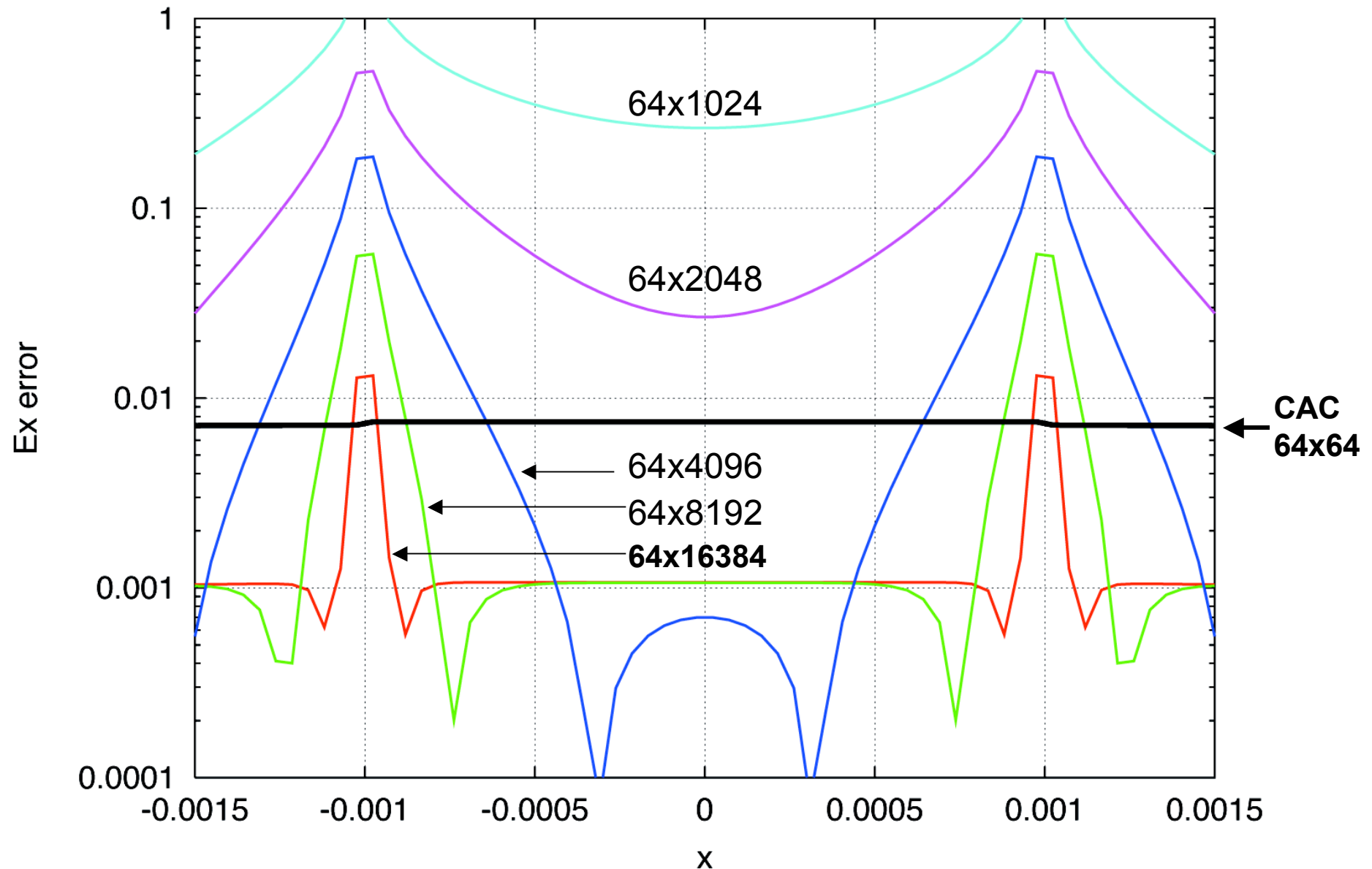
1:1000 test case; E_x vs. x : Standard Hockney Algorithm has huge errors



Ex vs. x : Reduced Vertical Scale



New algorithm (CAC) provides better than 1% accuracy using 64x64 grid (black curve). Old algorithm is worse everywhere on 64x2048 (purple); error > 10% at some locations on 64x4096 (blue); requires 64x16384 to achieve < 1% everywhere (red).



Bottom Line



- For this test problem, the standard Hockney algorithm would require ~500 times more computational effort to achieve the same worst-case accuracy as a simulation using the new ACE approach.
- CAC works whether the aspect ratio is large, small, or near unity, i.e. it is generally applicable.
- Verified in 2D; we expect the same to hold for 3D

Extension to Beams in Pipes



- CAC provides a crucial advantage, since the Green function falls off exponentially in z , though $\square(z)$ may change slowly over meters

- Due to shielding, sum can be truncated in the “long” direction:

$$\square_{i,j} = \sum_{i'=1}^{N_x} \sum_{j'=j-j_{cutoff}}^{j+j_{cutoff}} G_{i \square i', j \square j'}^{eff} \square_{i',j'}$$

- For long beam in a conducting pipe, if grid length in z is \gg pipe radius, can truncate at nearest neighbors:

$$\square_{i,j} = \sum_{i'=1}^{N_x} (G_{i \square i', j \square j-1}^{eff} \square_{i',j-1} + G_{i \square i', j}^{eff} \square_{i',j} + G_{i \square i', j+1}^{eff} \square_{i',j+1})$$

- For a rectangular pipe, can rewrite Green function as a sum of convolutions and correlations; then can still use FFT-based approach to sum over elements